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Micro-Drops and Charges: From Young to Lippmann and beyond

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ELECTROWETTING wetting enhancement



http://www.ee.duke.edu/Research/microfluidics/



THE BROADER CONTEXT

Roughness design towards reversible non- / full-wetting surfaces: From Fakir Droplets to Liquid Films

The 'HYDROFAKIR' project is approved for funding by the European Research Council, through an IDEAS Starting Grant awarded to Dr. Thanasis Papathanasiou. Duration: 5 years.





http://www.chemeng.ntua.gr/people/pathan/hydro.htm

• the transition from Cassie to Wenzel is NOT reversible

• a 'kick' is required to recover the Cassie state

AIM: design roughness + design 'kick' to make the transition reversible

IMPACT: self cleaning, tunable flow resistance, efficient liquid handling without mechanical parts, water collection from dew formation . . .



The HYDROFAKIR key parts: Design – Fabricate – Test





ELECTROWETTING IN A SESSILE DROP

Lippmann's equation:

Lippmann

200

Applied Voltage, Volt

250

300

350

60

40 20

50

100

150

Experiments

400



Limiting phenomena in electrowetting

- contact angle saturation
- droplet emission around the contact line

How electrowetting modifies the wettability ?



Kang, *Langmuir* **18**, 10318 (2002)

DROPLET ACTUATION



ELECTROWETTING ACTUATION

principle of operation

it works pretty nice !







APPLICATIONS

They will soon appear in consumer devices (microcameras in mobile phones, computer screens)

electrowetting display

liquid lens with variable focal length



Berge & Peseux, *European Physical Journal E* **3**, 159 (2000). Hayes & Feenstra, *Nature* **425**, 383 (2003).



BASIC QUESTIONS

What's beyond Young and Lippmann ?

- is the exact solution of the electro-hydrostatic problem enough for explaining the contact angle saturation ?
- if no, how can we use it to extract information and guide further analysis?



ELECTRO-HYDROSTATICS

governing equations and boundary conditions



Young-Laplace equation $-g\delta\rho h + \varepsilon_0 E^2 / 2 + \gamma_{1v} C = K$ Field distribution Electric potential: $E \equiv \nabla u$ $\nabla^2 u = 0$ In the dielectric, in the air Volume constraint $\int_{0}^{\pi/2} f^3 \sin\theta d\theta = 2$

BOUNDARY CONDITIONS:

FIELD CONDITIONS - ASYMPTOTIC CONDITIONS - INTERFACIAL CONDITIONS



THE COMPUTATIONAL PROBLEM



THE EQUATIONS ARE SOLVED FOR f, u SIMULTANEOUSLY



WHY SATURATION?

- H. J. J. Verheijen and M. W. J. Prins, *Langmuir* 15, 6616 (1999).
 trapping of charge
- M. Vallet, B. Berge and L. Vovelle, *Polymer* 37, 2465 (1996).
 air ionization at the contact line
- E. Seyrat and R. A. Hayes, J. Appl. Phys. 90, 1383 (2001). increasing breakdown strength, improves electrowetting performance
- V. Peykov, A. Quinn and J. Ralston, *Colloid Polym. Sci.* 278, 789 (2000).
- A. Quinn, R. Sedev and J. Ralston, J. Phys. Chem. B 109, 6268 (2005).
 saturation when E_{sl} = 0
- B. Shapiro et al., *J. Appl. Phys.* 93, 5794 (2003).
 relative conductivity of the media is important
- A. G. Papathanasiou and A. G. Boudouvis, *Appl. Phys. Lett.* 86, 164102 (2005). breakdown strength determines the onset of the saturation



THE COMPUTED DROP SHAPE



15

d = 1 μ m SiO₂, ϵ_r = 3.8 γ_{lv} =0.072 N/m, water θ_Y = 120° θ_V = 70° at V = 75V



SHAPE DEPENDENCE ON THE APPLIED VOLTAGE





PREDICTING THE ONSET OF THE SATURATION



$$\label{eq:eq:constraint} \begin{split} d &= 1 \; \mu m \; SiO_2, \; \epsilon_r = 3.8 \\ E_{bd} &= 10 \; MV \; / \; cm \\ \gamma_{Iv} &= 0.072 \; N/m, \; water \\ \theta_{sat} &\sim 75^\circ \; at \; V \; \sim \; 75 \; V \end{split}$$

Papathanasiou & Boudouvis, Applied Physics Letters 86, 164102 (2005).



More findings More questions

If the main cause of the saturation is the local breakdown of the dielectric at the three phase contact line then :

a) how can we show it experimentally ?

b) how can we improve our computational analysis to predict the asymptotic stabilization of the contact angle past the saturation?

c) what can we do to inhibit the saturation and improve electrowetting?



EXPERIMENTS

(the setup)





LEAKAGE CURRENT MEASUREMENTS



if the contact angle saturation is connected with local dielectric breakdown then significant increase of the leakage current is expected !

d = 130 nm SiO₂, ϵ_r = 3.8 γ_{lv} =0.072 N/m, water θ_Y = 117°

Papathanasiou, Papaioannou & Boudouvis, Journal of Applied Physics 103, 034901 (2008).



Is the breakdown localized at the contact line ?



d = 500 nm, TEOS, $\epsilon_r = 3.8$ $\gamma_{lv}=0.072 \text{ N/m, water}$ $\theta_Y = 117^{\circ}$



MULTILAYER vs SINGLE-LAYER DIELECTRICS

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 $E_{bd, ONO} > E_{bd, SiO_2}$

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Holland, IEEE Trans. Nucl. Sci. 42, 423 (1995).

Nozaki & Giridhar, IEEE Electron Device Lett. 7, 486 (1986).



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IMPROVING THEORETICAL PREDICTIONS

- the field strength cannot exceed the breakdwon strength of the solid dielectric
- regions with higher field strength are switched to conductive





SIMULATING ELECTROWETTING BEYOND THE SATURATION



d =1000 nm, SiO₂, ϵ_r = 3.8, γ_{lv} =0.072 N/m, water, θ_{γ} = 117°

Drygiannakis, Papathanasiou & Boudouvis, Langmuir 25 (2009).



Potential and field strength distribution at the TPL





TRAPPED CHARGE DISTRIBUTION



 $\label{eq:constraint} \begin{array}{l} d = 1 \ \mu m \ SiO_2, \ \epsilon_r = 3.8 \\ E_{bd} = 10 \ MV \ / \ cm \\ \gamma_{lv} = 0.072 \ N/m, \ water \end{array}$



TESTING PREDICTIONS FOR DIFFERENT MATERIALS



Parylene and PTFE measurements from: Verheijen & Prins, *Langmuir* **15**, 6616 (1999) Vallet, Vallade & Berge, *European Physical Journal B* **11**, 583 (1999)



CONCLUSIONS

breakdown strength – charge trapping – contact angle saturation

succesfull simulation / prediction: onset of the saturation asymptotic stabilization of the contact angle trapped charge (naturally computed)

based only on material properties

the saturation is evidently shifted to higher voltages when composite dielectrics with higher breakdown strength are used



Collaborators / students

- Dr. Thanasis Papathanasiou
- Dr. Antonis Drygiannakis
- George Pashos, Doctoral student

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