

Modeling and Simulation of Viscoelastic Flows in Polymer Processing with Integral Constitutive Equations

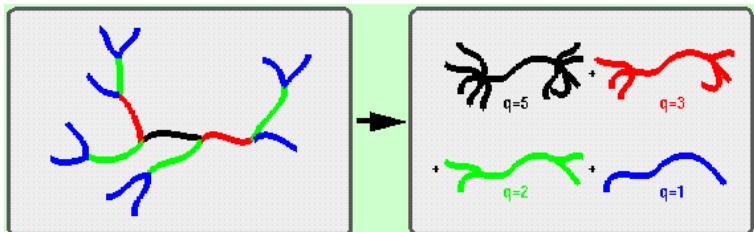
Evan Mitsoulis

CAMP - R&D

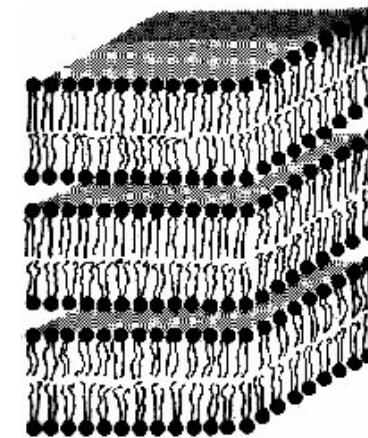
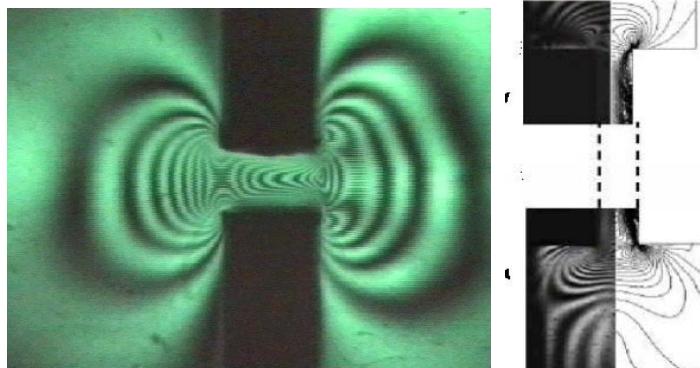
Computer-Aided Materials Processing – Rheology & Design
School of Mining Engineering and Metallurgy
National Technical University of Athens
Athens, Greece



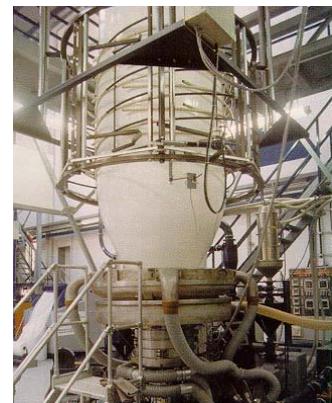
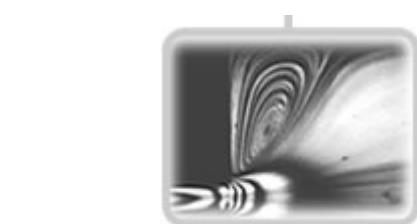
Molecular Structure



Processing



Product (Structure-Properties)



A. GOVERNING FLOW EQUATIONS (Conservation equations)

A1. Equation of conservation of mass

$$\nabla \cdot \bar{v} = 0$$

A2. Equation of conservation of momentum

$$\rho \bar{v} \cdot \nabla \bar{v} = -\nabla p + \nabla \cdot \bar{\tau} + \rho \bar{g}$$

A3. Equation of conservation of energy

$$\rho C_p \bar{v} \cdot \nabla T = k \nabla^2 T + \bar{\tau} : \nabla \bar{v}$$

B. RHEOLOGICAL MODEL (Constitutive equation)

$$\bar{\tau} = f(\bar{v}, \nabla \bar{v})$$



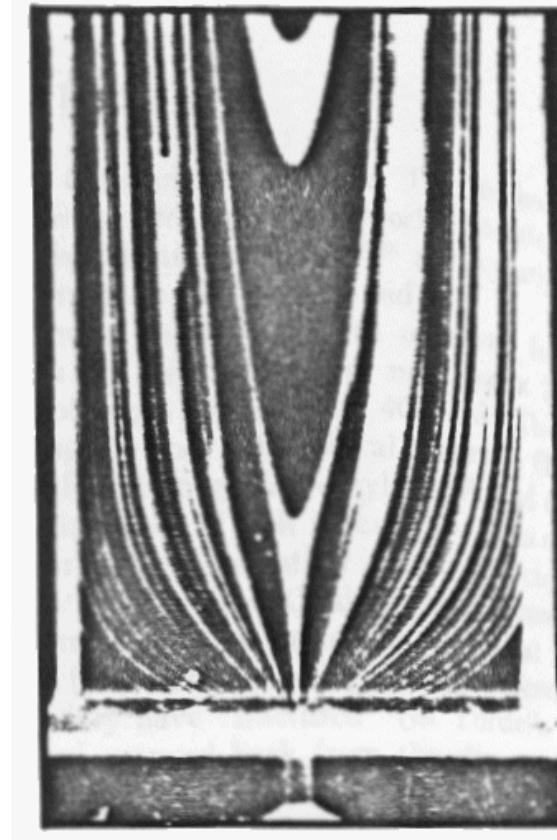
I Think
It's **STRESS**



20:1 Axisymmetric Contraction

LDPE HDPE

$$\tau_w = 0.1 \text{ MPa}$$



Bagley and Birks, J. Appl. Phys., 1960



INTEGRAL CONSTITUTIVE EQUATIONS

$$\bar{\bar{\tau}}(t) = \int_{-\infty}^t [M_1(t-t', I_{C^{-1}}, I_C) \bar{\bar{C}}_t^{-1}(t') + M_2(t-t', I_{C^{-1}}, I_C) \bar{\bar{C}}_t(t')] dt'$$

where:

M_1, M_2 = material – property functions

t = present time

t' = past time

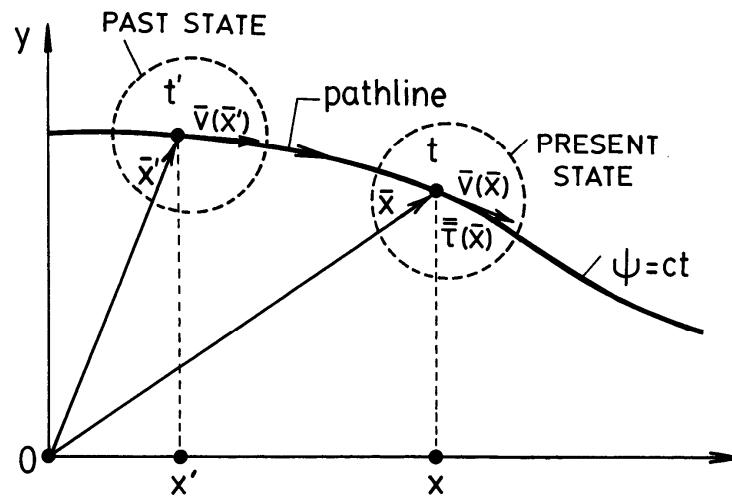
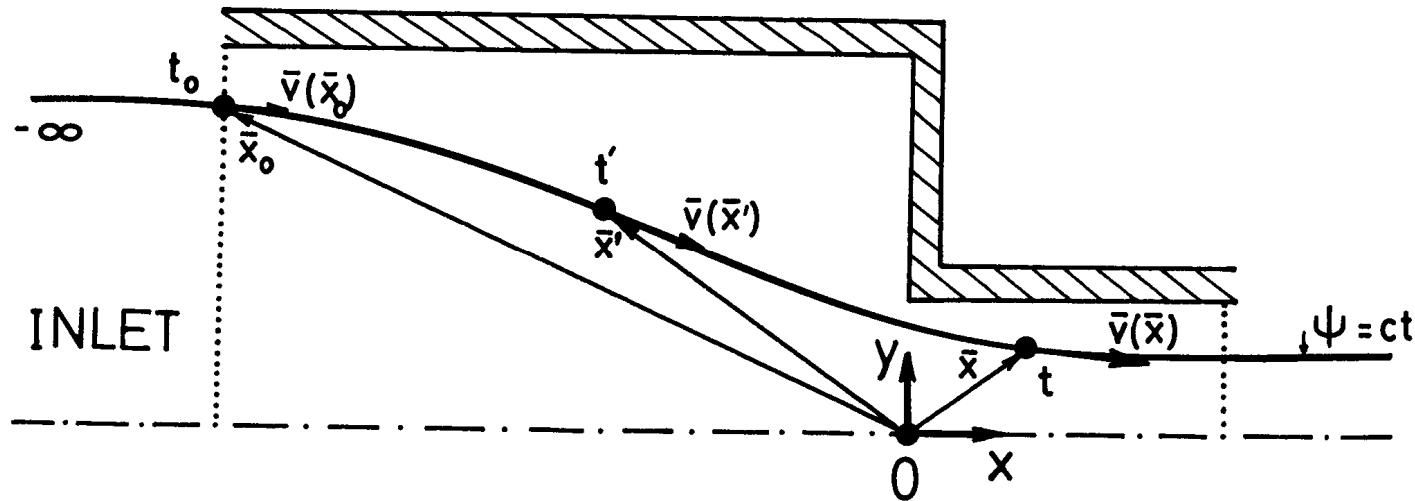
$I_{C^{-1}}, I_C$ = first invariants

$\bar{\bar{C}}_t$ = Cauchy – Green tensor

$\bar{\bar{C}}_t^{-1}$ = Finger strain tensor



Particle Tracking



INTEGRAL CONSTITUTIVE EQUATIONS

K-BKZ EQUATION

- Kaye (1962), Bernstein-Kearsley-Zapas (1963)

$$\bar{\bar{\tau}}(t) = \int_{-\infty}^t \left[\frac{\partial U(t-t', I_{C^{-1}}, I_C)}{\partial I_{C^{-1}}} \bar{\bar{C}}_t^{-1}(t') - \frac{\partial U(t-t', I_{C^{-1}}, I_C)}{\partial I_C} \bar{\bar{C}}_t(t') \right] dt'$$

- Factorized K-BKZ proposed by White and Tokita (1967)

$$\bar{\bar{\tau}}(t) = \int_{-\infty}^t m(t-t') \left[\frac{\partial V(I_{C^{-1}}, I_C)}{\partial I_{C^{-1}}} \bar{\bar{C}}_t^{-1}(t') - \frac{\partial V(I_{C^{-1}}, I_C)}{\partial I_C} \bar{\bar{C}}_t(t') \right] dt'$$

where: U = potential function

$I_{C^{-1}}, I_C$ = first invariants

$m(t-t')$ = time-dependent function

$\bar{\bar{C}}_t$ = Cauchy-Green tensor

V = strain-dependent function

$\bar{\bar{C}}_t^{-1}$ = Finger strain tensor



The Faces behind the Names

4.14. BKZ



Barry Bernstein, Elliot Kearsley and Louis Zapas at the Rochester Society of Rheology meeting in 1991.

**From Rheology: An Historical Perspective,
R.I. Tanner & K. Walters, Elsevier, 1998.**



INTEGRAL CONSTITUTIVE EQUATIONS

RIVLIN-SAWYERS EQUATION

- Rivlin and Sawyers (1971)

$$\bar{\bar{\tau}}(t) = \int_{-\infty}^t \left[W_1(t-t', I_{C^{-1}}, I_C) \bar{\bar{C}}_t^{-1}(t') + W_2(t-t', I_{C^{-1}}, I_C) \bar{\bar{C}}_t(t') \right] dt'$$

- Factorized R-S

$$\bar{\bar{\tau}}(t) = \int_{-\infty}^t m(t-t') \left[H_1(I_{C^{-1}}, I_C) \bar{\bar{C}}_t^{-1}(t') + H_2(I_{C^{-1}}, I_C) \bar{\bar{C}}_t(t') \right] dt'$$

where: W = not the potential function

$I_{C^{-1}}, I_C$ = first invariants

$m(t-t')$ = time-dependent function

$\bar{\bar{C}}_t$ = Cauchy - Green tensor

H = strain-dependent function

$\bar{\bar{C}}_t^{-1}$ = Finger strain tensor



INTEGRAL CONSTITUTIVE EQUATIONS

Special Cases of the Factorized RIVLIN-SAWYERS EQUATION

- Wagner Model (1980)

$$M = \left[\sum_k \frac{a_k}{\lambda_k} \exp\left(-\frac{t - t'}{\lambda_k}\right) \right], \quad H = \exp\left(-\beta \sqrt{\alpha I_{C^{-1}} + (1 - \alpha) I_C - 3}\right)$$

- Papanastasiou-Scriven-Macosko (PSM) Model (1983)

$$M = \left[\sum_k \frac{a_k}{\lambda_k} \exp\left(-\frac{t - t'}{\lambda_k}\right) \right], \quad H = \left[\frac{\alpha}{(\alpha - 3) + \beta I_{C^{-1}} + (1 - \beta) I_C} \right]$$

where:

λ_k = relaxation time

$I_{C^{-1}}, I_C$ = first invariants

a_k = relaxation modulus

$\overline{\overline{C}}_t$ = Cauchy – Green tensor

α, β = material constants

$\overline{\overline{C}}_t^{-1}$ = Finger strain tensor



Relevant Literature

- **Kaye**, A. (1962) Non-Newtonian Flow in Incompressible Fluids, Part I: A General Rheological Equation of State; Part II: Some Problems in Steady Flow, Note No. 134, College of Aeronautics, Cranford, UK.
- **Bernstein**, B., **Kearsley**, E.A., and **Zapas**, L. (1963) A Study of Stress Relaxations with Finite Strain, *Trans. Soc. Rheol.*, **7**, 391-410.
- **Papanastasiou**, A.C., **Scriven**, L.E., and **Macosko**, C.W. (1983) An Integral Constitutive Equation for Mixed Flows: Viscoelastic Characterization, *J. Rheol.* **27**, 387-410.
- **Olley**, P. (2000) An adaptation of the separable KBKZ equation for comparable response in planar and axisymmetric flow, *J. Non-Newtonian Fluid Mech.*, **95**, 35-53.



Numerical Methods (2-D)

- ◆ u-v-p formulation (primitive variables)
- ◆ u-v-p formulation (SI)
- ◆ mixed formulation (u-v-p- τ) (EVSS)
- ◆ enhanced mixed formulation (u-v-p- τ -g) (DEVSS-G)
- ◆ many others (EEME, SI, EVSS/SUPG, EVSS/SU,
AVSS/SI, AVSS/SUPG, etc.)



Numerical Methods (3-D)

- ◆ 3-D x-y-z-p-t formulation (primitive variables)
- ◆ Lagrangian Integral Method (LIM)
- ◆ H.K. Rasmussen, “Time-dependent finite-element method for the simulation of three-dimensional viscoelastic flow with integral models”, JNNFM, 84, 217-232 (1999).

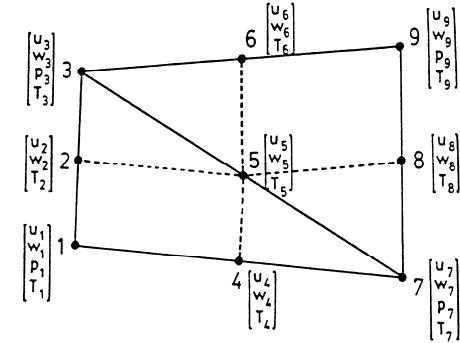


Relative Difficulty of Pom-Pom vs. K-BKZ Models (or differential vs. integral models)

2-D dof / element

K-BKZ

Pom-pom



8 u (serendipity)

8 v

4 p

8 T

20 (28)

planar

$$4 (\tau_{xx}, \tau_{yy}, \tau_{xy}) = 12/\text{mode} \times 8 = 96$$

$$4 (g_{xx}, g_{yy}, g_{xy}) = \underline{12/\text{mode} \times 8 = 96}$$

46 (55)

214 (223)

axisymmetric

$$4 (\tau_{rr}, \tau_{zz}, \tau_{rz}, \tau_{\theta\theta}) = 16/\text{mode} \times 8 = 128$$

$$4 (g_{rr}, g_{zz}, g_{rz}, g_{\theta\theta}) = \underline{16/\text{mode} \times 8 = 128}$$

54 (63)

278 (287)

dof/elem/8 modes



Parameters for Fitting Data for the IUPAC-LDPE Melt

Pom-Pom

k	λ_k (s)	a_k (Pa)	q_k	$(\tau_b/\tau_s)_k$
1	0.0001	1.29×10^5	1	4.5
2	0.001	9.48×10^4	1	3.9
3	0.01	5.86×10^4	2	3.7
4	0.1	2.67×10^4	2	3.6
5	1.0	9.80×10^3	3	2.9
6	10.0	1.89×10^3	7	2.0
7	100.0	1.80×10^2	8	2.1
8	1000.0	1.00×10^0	9	1.3

Inkson, McLeish, Harlen, Groves, *J. Rheol.* (1999)

K-BKZ (PSM)

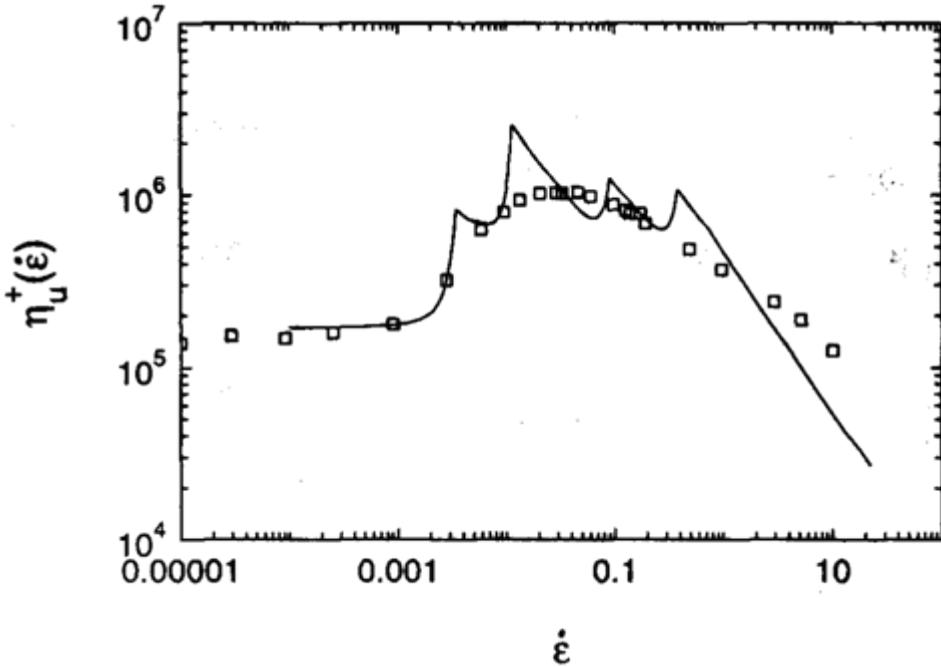
k	λ_k (s)	a_k (Pa)	α_k	β_k
1	0.0001	1.29×10^5	14.38	0.018
2	0.001	9.48×10^4	14.38	0.018
3	0.01	5.86×10^4	14.38	0.08
4	0.1	2.67×10^4	14.38	0.12
5	1.0	9.80×10^3	14.38	0.12
6	10.0	1.89×10^3	14.38	0.16
7	100.0	1.80×10^2	14.38	0.03
8	1000.0	1.00×10^0	14.38	0.002

Barakos and Mitsoulis, *J. Rheol.* (1995)

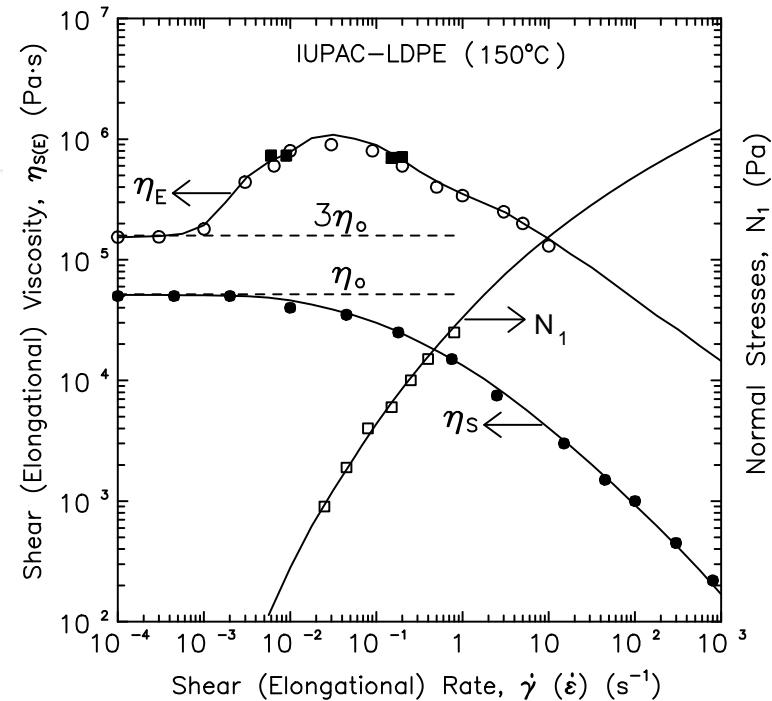


Elongational Viscosity

Pom-Pom



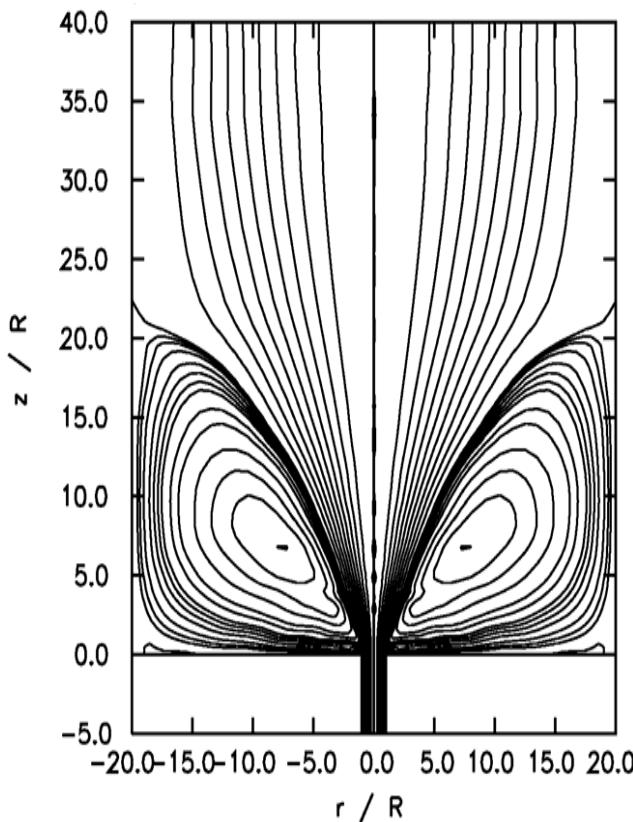
K-BKZ (PSM)



LDPE Melts

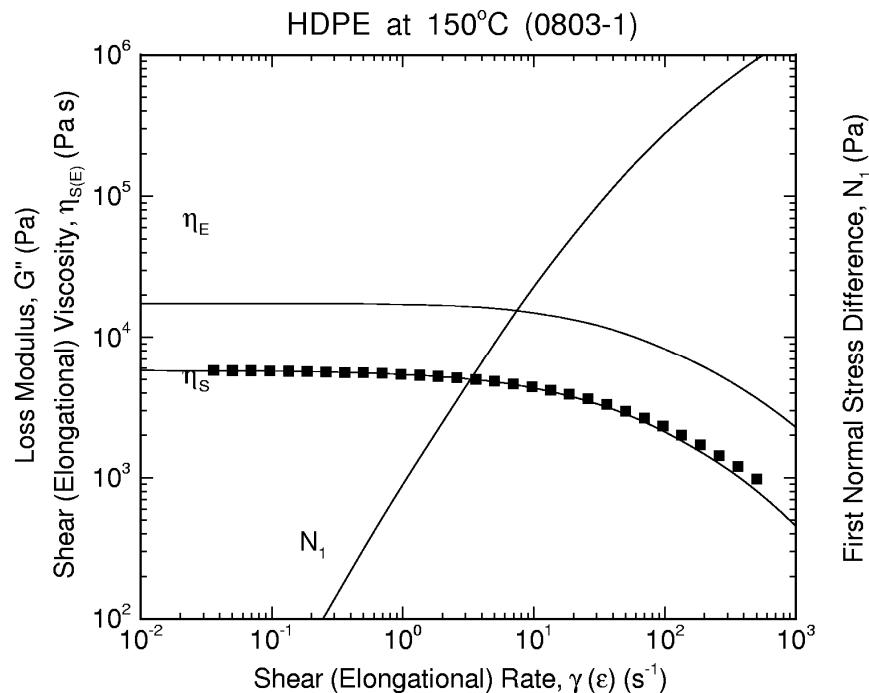
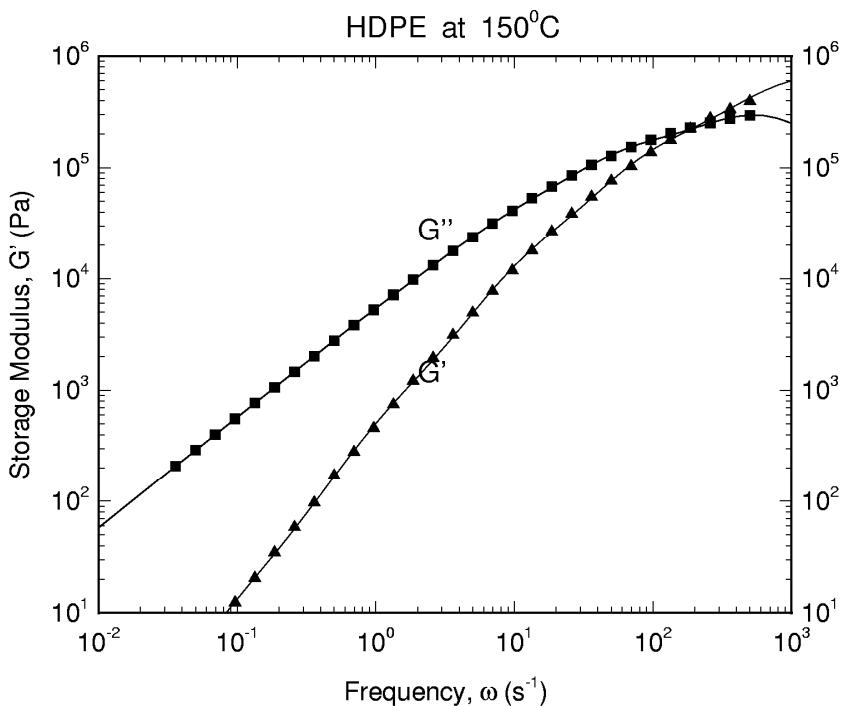
Simulations Experiments

$$\gamma_a = 84 \text{ s}^{-1}, \tau_w = 0.1 \text{ MPa}$$



HDPE Melt (0803-1 at 150°C)

K-BKZ Model Original Fit



- ◆ A non-linear regression analysis is performed on the K-BKZ integral constitutive equation to determine the parameters in the model.



HDPE Melt (0803-1 at 150°C)

K-BKZ Model Parameters

Original Fit

k	λ_k (s)	a_k (Pa)	α_k	β_k
1	0.001723	5.46×10^5	4.0	0.9985
2	0.01299	1.73×10^5	4.0	0.9985
3	0.08644	2.15×10^4	4.0	0.9985
4	0.7991	7.69×10^2	4.0	0.9985
5	9.398	1.39×10^1	4.0	0.9985

$$\theta = 0$$



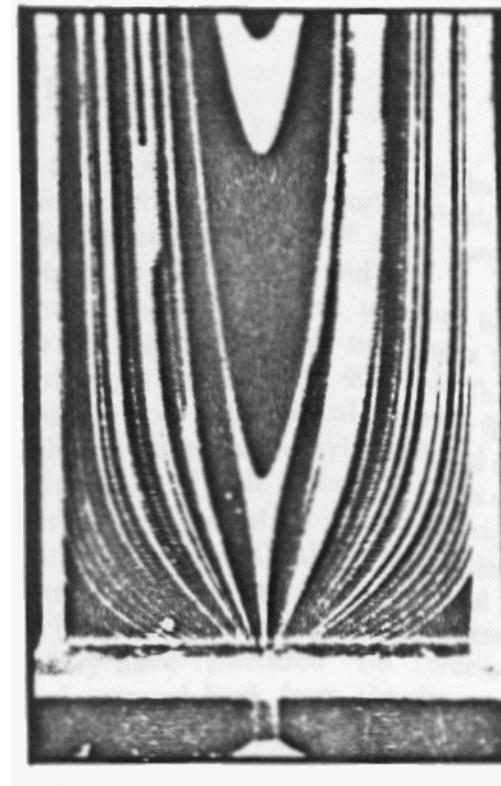
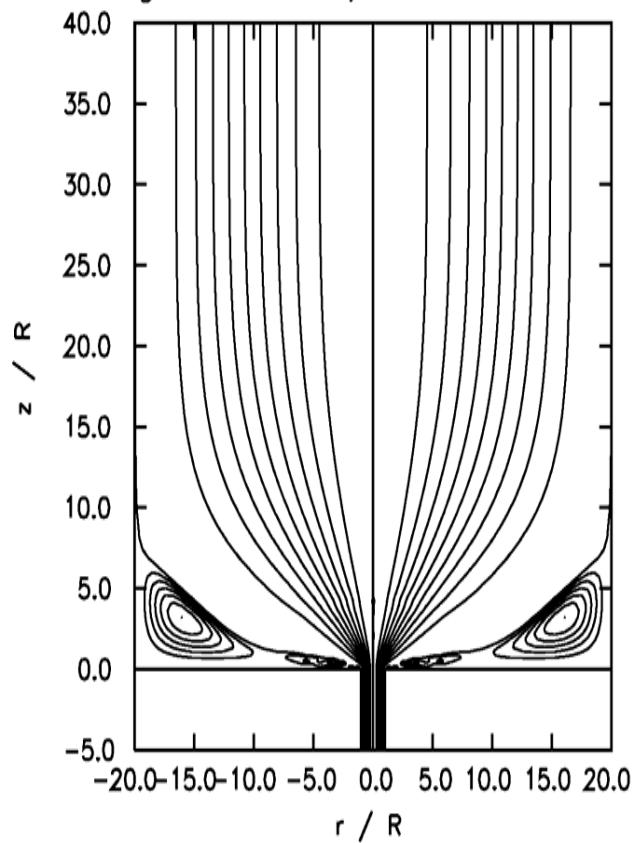
P. Wood-Adams, PhD Thesis, McGill University, 1999.



HDPE Melts

Simulations Experiments

$$\gamma_a = 248 \text{ s}^{-1}, \tau_w = 0.1 \text{ MPa}$$



NON-ISOTHERMAL CONSTITUTIVE EQUATION

TIME-TEMPERATURE SHIFTING CONCEPT

$$\bar{\tau}(t) = \frac{1}{1-\theta} \int_{-\infty}^{\xi(t)} M H \left[C_{t(\xi)}^{-1} (t'(\xi')) + \theta C_{t(\xi)} (t'(\xi')) \right] d\xi'$$

$$M = \left[\sum_k \frac{a_k}{\lambda_k} \exp\left(-\frac{\xi - \xi'}{\lambda_k}\right) \right], \quad H = \left[\frac{\alpha}{(\alpha - 3) + \beta I_{C^{-1}} + (1 - \beta) I_C} \right]$$

where:

λ_k = relaxation time

a_k = relaxation modulus

α, β, θ = material constants

$I_{C^{-1}}, I_C$ = first invariants

\bar{C}_t = Cauchy – Green tensor

\bar{C}_t^{-1} = Finger strain tensor

Note: $\frac{N_2}{N_1} = \frac{\theta}{1-\theta}$



Non-Isothermal Analysis

- Peclet Number:

$$Pe = \frac{\rho C_p \bar{V} R}{k}$$

- Nahme Number:

$$Na = \frac{a \eta_o \bar{V}^2}{k}$$

where

or

$$a(T) = \frac{E_o}{R_g T_o^2}$$

$$a(T) = \frac{\eta}{\eta_o} = \exp \left[-\frac{E_o}{R_g} \left(\frac{1}{T} - \frac{1}{T_o} \right) \right]$$

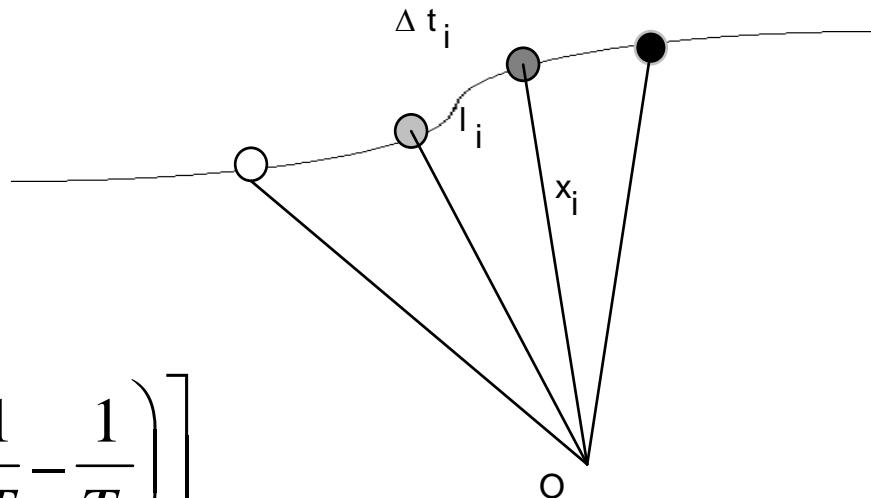


Pseudo-Time Integral Method

$$d\xi = \frac{dt'}{a(T(t'))}$$

$$dt' = \frac{dl}{dV}$$

$$a(T) = \frac{\eta}{\eta_0} = \exp\left[-\frac{E_0}{R_0}\left(\frac{1}{T} - \frac{1}{T_0}\right)\right]$$

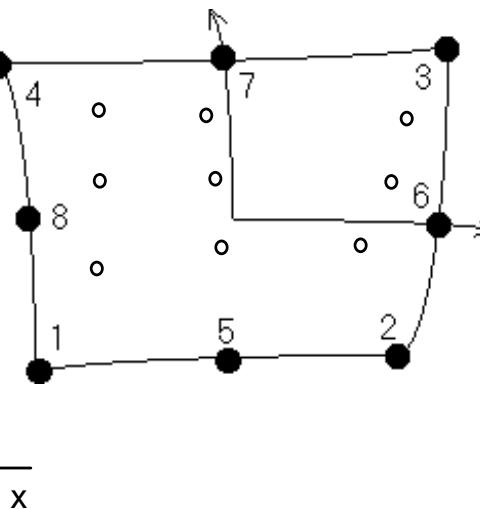
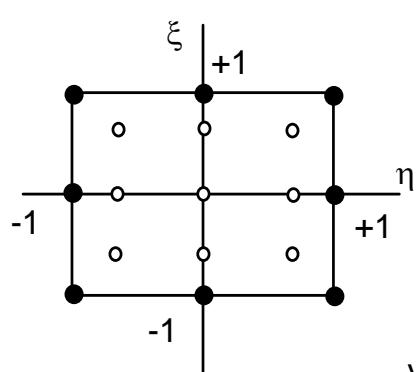


t' : observer's time

ξ : particle's time



'UPWIND' FEM



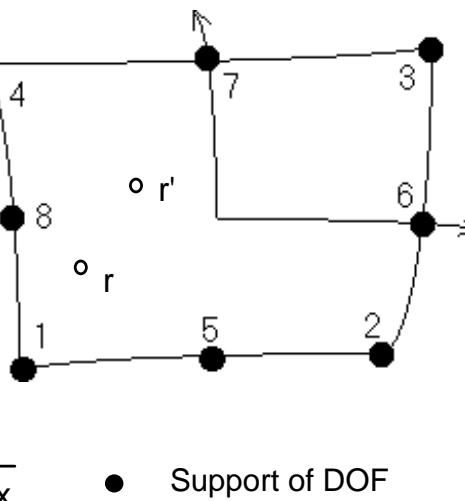
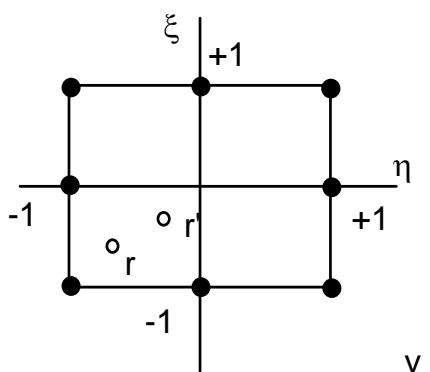
$$r' - r = a \frac{h}{2}$$

$$a = \coth\left(\frac{Pe}{2}\right)$$

$$Pe = \frac{|u| h}{k}$$

h =length of streamline

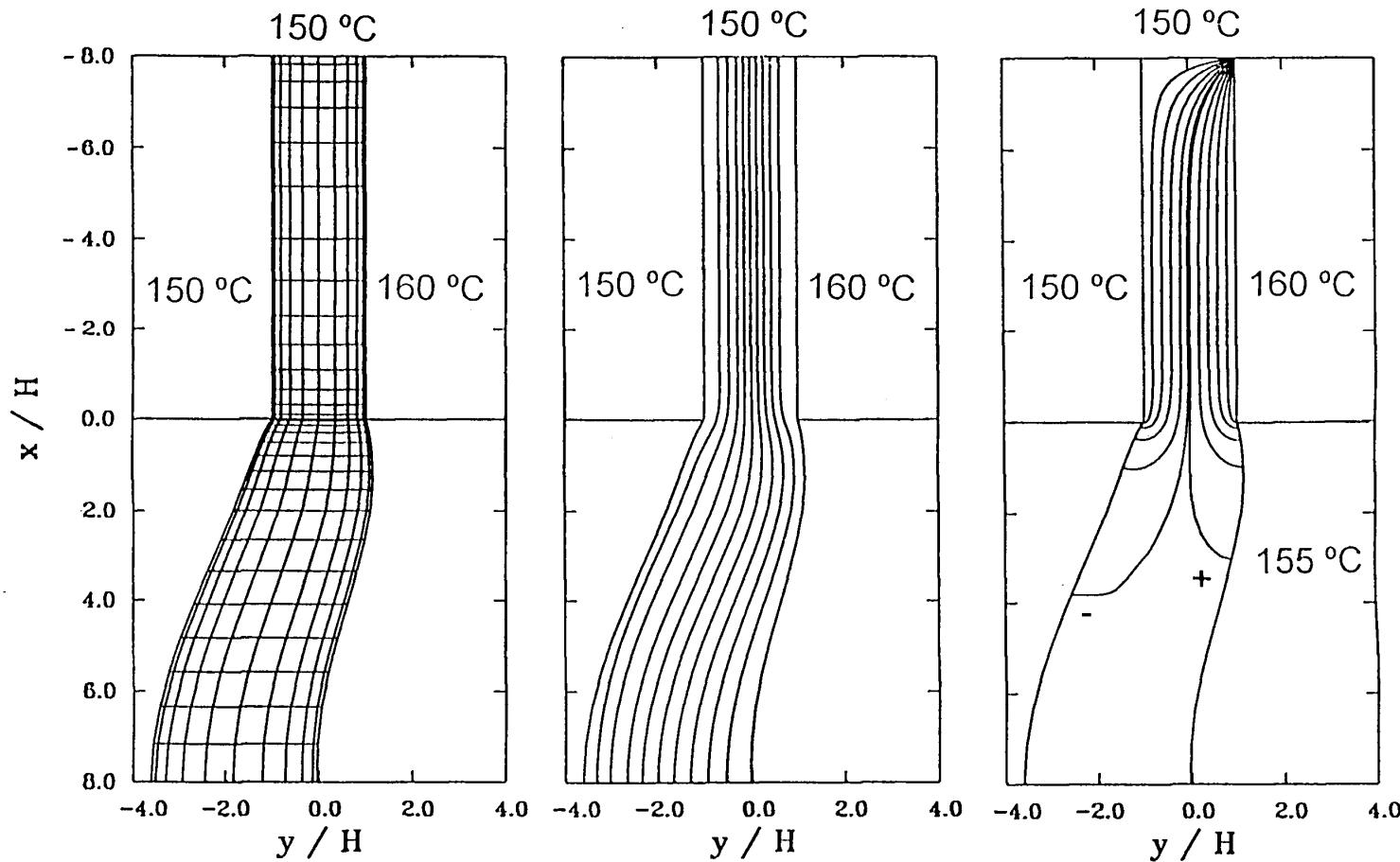
k =diffusion coefficient



- Support of DOF
- Support of Integration Node



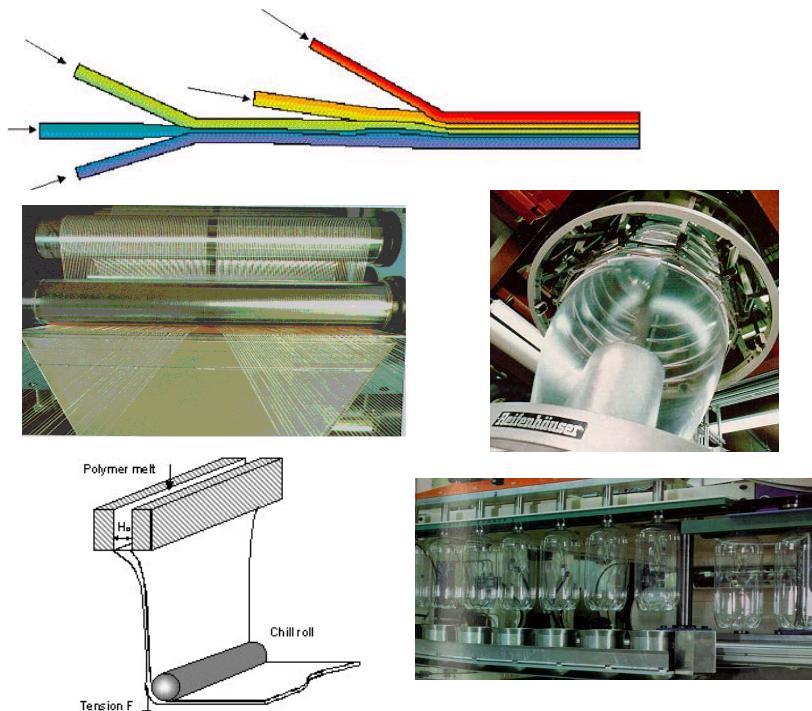
Extrudate Swell and Bending Phenomenon



Polymer Processes

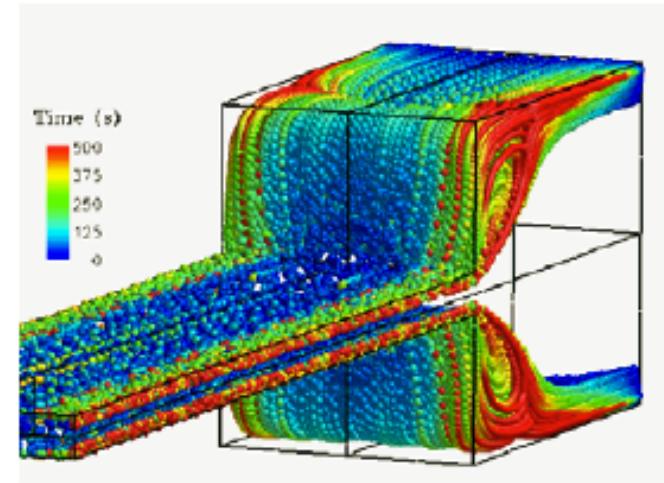
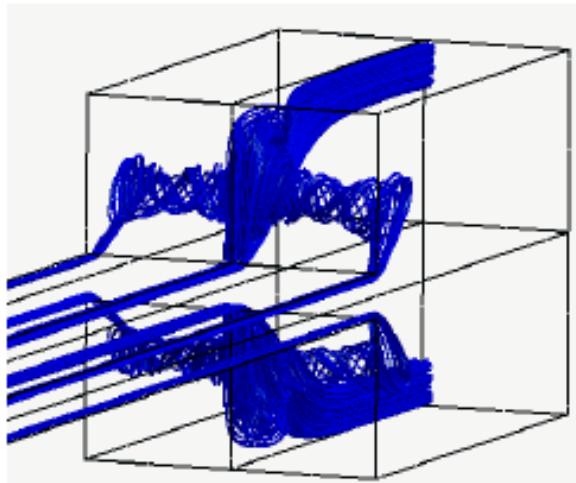
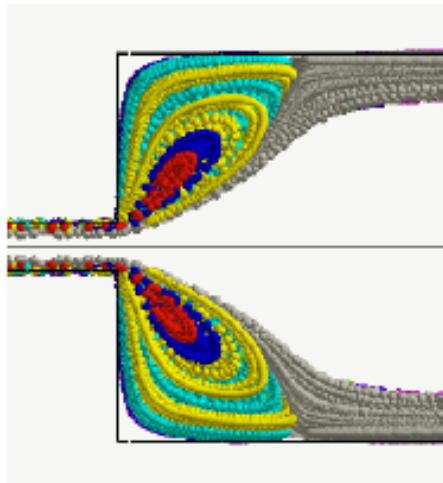
Polymer processes where the K-BKZ model has been used:

- Coextrusion
- Fiber Spinning
- Film Blowing
- Film Casting
- Blow Molding



New Developments 3-D Simulations

Pom-Pom vs K-BKZ



CONCLUSIONS

- Computational rheology has made good progress in the last 15 years due to rapid growth in computing power.
- From the modelling point of view, progress has been achieved mainly by using integral constitutive equations of the K-BKZ type, which are quite capable of fitting rheological data for polymer solutions and melts in simple flows.
- In the processing of non-Newtonian, complex materials, many unsolved problems still remain; however there are now available appropriate means of tackling them, so that the analysis and design of several processes becomes feasible.



FURTHER WORK

- The subject of **viscoelasticity** requires further study for more appropriate **damping functions** for fast deformations and better (faster) computational methods.
 - **Time-dependent flows**
 - **Multiple-layer flows**
 - **Three-Dimensional flows**

